

Communication

Azafullerene Encapsulated Single-Walled Carbon Nanotubes with n-Type Electrical Transport Property

Toshiro Kaneko, Yongfeng Li, Shohei Nishigaki, and Rikizo Hatakeyama

J. Am. Chem. Soc., 2008, 130 (9), 2714-2715 • DOI: 10.1021/ja0773960

Downloaded from http://pubs.acs.org on February 8, 2009



More About This Article

Additional resources and features associated with this article are available within the HTML version:

- Supporting Information
- Links to the 1 articles that cite this article, as of the time of this article download
- Access to high resolution figures
- Links to articles and content related to this article
- Copyright permission to reproduce figures and/or text from this article

View the Full Text HTML





Published on Web 02/08/2008

Azafullerene Encapsulated Single-Walled Carbon Nanotubes with n-Type Electrical Transport Property

Toshiro Kaneko,* Yongfeng Li, Shohei Nishigaki, and Rikizo Hatakeyama

Department of Electronic Engineering, Tohoku University, Sendai 980-8579, Japan

Received September 25, 2007; E-mail: kaneko@ecei.tohoku.ac.jp

The azafullerene $C_{59}N$, i.e., a single carbon atom of the fullerene C₆₀ cage replaced by a nitrogen atom,¹⁻³ has attracted special attention because of its interesting properties and applications in superconductivity, photoelectric devices, and organic semiconductors.⁴ Although the theoretical study has indicated that the azafullerene has the property of an electron donor,⁵ to our knowledge, there is no experimental verification of it. On the other hand, following the discovery of carbon nanotubes, carbon peapods, i.e., single-walled carbon nanotubes (SWNTs) encapsulating fullerenes, are recognized to be promising materials as potential building blocks for nanoelectronics. The presence of fullerenes inside SWNTs is expected to significantly modify the band structure of SWNTs and consequently affect their electrical transport properties. It has been reported experimentally that fullerene peapods can exhibit different transport properties, such as p-type, ambipolar, and metallic behaviors.⁶ However, an n-type characteristic of the fullerene peapod has never been reported, which is indispensable for realizing SWNT-based electronic devices.

In this communication, we investigate synthesis, characterization, and electric transport properties of SWNTs encapsulating $C_{59}N$, which results in a clear-cut observation of the n-type behavior of the peapod. It is the first time that the azafullerene $C_{59}N$ is proven to be a good electron donor.

The azafullerene C₅₉N is synthesized using a plasma irradiation method under the following conditions: plasma density $n_{\rm p} \approx 10^9$ ${\rm cm}^{-3}$, electron temperature $T_{\rm e} \approx 0.5 {\rm eV}$, and nitrogen-ion irradiation energy $E_i = 10-40 \text{ eV}$.⁷ The fullerene C₆₀ after plasma irradiation is dissolved in toluene, and its mixture is separated into a residue and a solution. The mass spectroscopy analysis of the formed azafullerene is performed using a laser-desorption time-of-flight mass spectrometer (LD-TOF-MS, Shimadzu AXIMA-CFR+). Figure 1 shows mass spectra of (a) the residue and (b) the solution of the fullerene dissolved toluene in a positive ion mode of LD-TOF-MS. The peak at the mass number 722 is the most distinct in the residue, which originates from azafullerene C₅₉N. It is known that $C_{59}N$ is stable in the form of $(C_{59}N)_2$, but the mass number of $C_{59}N$ monomer (722) is detected here because the chemical bonding between C₅₉N monomers is broken by the laser irradiation. In the solution, however, the peak intensity corresponding to C₅₉N is much less than that of C_{60} (720) because $C_{59}N$ dissolves little in toluene. Therefore, C₅₉N is dominant in the residue, although the quantitative abundance ratio of C59N to C60 is not determined from the results of LD-TOF-MS.

This C₅₉N mixed with slight C₆₀ is encapsulated into SWNTs⁸ by either a vapor reaction method or a plasma ion-irradiation method. For the vapor reaction method, the purified SWNTs together with C₅₉N fullerene powders, are sealed in a glass tube under a vacuum condition of $\sim 10^{-5}$ Torr. After that, the sealed glass tube is heated at 420 °C for 48 h to encapsulate C₅₉N in SWNTs. Raw samples are obtained after the above process and then purified via a washing process in toluene to remove most of



Figure 1. Mass spectra of (a) the residue and (b) the solution of the fullerene dissolved toluene in positive ion mode of LD-TOF-MS.



Figure 2. Raman spectra for the pristine, C_{60} encapsulated, and $C_{59}N$ encapsulated SWNTs in the ranges (a) 100–400 cm⁻¹ (Radial breathing mode) and (b) 1100–1700 cm⁻¹. The Raman spectra are obtained at an excitation wavelength of 488 nm.

the excess fullerenes attached to the surface of SWNTs. The $C_{59}N$ encapsulated SWNTs purified are examined in detail by field emission transmission electron microscopy (FE-TEM, Hitachi HF-2000) operated at 200 kV and Raman spectroscopy (Jovin Yvon T-64000) with an Ar laser at 488 nm. The electronic transport properties of various SWNTs are investigated by fabricating them as the channels of field effect transistor (FET) devices. These SWNTs samples are ultrasonically dispersed in *N*,*N*-dimethylformamide first and then spincoated on FET substrates, each of which consists of Au drain–source electrodes on a SiO₂ insulating layer. A heavily doped Si substrate serves as a backgate. The detailed fabrication process for FET devices can be found elsewhere.^{9,10} The transport measurements are performed at room temperature under vacuum conditions on a semiconductor parameter analyzer (Agilent 4155C).

Figure 2 shows Raman spectra for the pristine, C_{60} encapsulated (C_{60} @SWNTs), and C_{59} N encapsulated SWNTs (C_{59} N@SWNTs) in the ranges (a) 100–400 cm⁻¹ (Radial breathing mode) and (b) 1100–1700 cm⁻¹. The pristine SWNTs show Raman peaks at 161 and 178 cm⁻¹, which correspond to the SWNTs with diameters of about 1.54 and 1.40 nm, respectively. The spectrum shape in the C_{60} or C_{59} N encapsulated SWNTs drastically changes in comparison with the pristine SWNTs, which is especially reflected in a decrease



Figure 3. TEM images for (a) C₆₀@SWNTs and (b) C₅₉N@SWNTs (scale bar is 2 nm).



Figure 4. Source-drain current (IDS) vs gate voltage (VG) characteristics at room temperature for (a) pristine SWNT, (b) C₆₀ encapsulated, and (c) C₅₉N encapsulated SWNTs.

in the peak intensity at 161 cm⁻¹. These results of Raman spectra in the RBM region give indirect evidence of the encapsulation of other materials, i.e., C₆₀ or C₅₉N, inside SWNTs. In addition, it is found that the frequency of the G⁺ mode (1591 cm⁻¹) slightly upshifts and down-shifts in the cases of C60 and C59N encapsulations, respectively (see Supporting Information).

To further verify the encapsulation of the fullerenes, Raman spectra in the range 1100-1700 cm⁻¹ are measured in detail as shown in Figure 2b. The clear peaks at 1423 and 1469 cm⁻¹ are observed in the C_{60} @SWNTs, corresponding to the modes Hg(7) and Ag(2) for C₆₀, respectively.¹¹ However, these peaks disappear in the C₅₉N@SWNTs because the peak intensities of Raman spectra of the modes Hg(7) and Ag(2) for C₅₉N are very weak compared with those for C_{60} (see Supporting Information).

A TEM observation result of C₆₀@SWNTs is shown in Figure 3a, which indicates that C₆₀ molecules with spherical symmetry are filled into SWNTs, where C₆₀ is isolated from each other. In the case of C₅₉N@SWNTs, on the other hand, paired C₅₉N dimers which are indicated by arrows in Figure 3b are observed along the tube axis. This means that C59N is actually encapsulated into SWNTs and makes the dimer, forming a one-dimensional chainlike structure inside the SWNTs. The TEM images for C₅₉N encapsulated SWNTs with lower magnification show the relatively high filling rate of the fullerenes (see Supporting Information).

The transport property of pristine semiconducting SWNTs is well-known to exhibit the p-type behavior as shown in Figure 4a, where a characteristic curve of source-drain current I_{DS} versus gate voltage $V_{\rm G}$ is described for source-drain voltage $V_{\rm DS} = 1$ V. Figure 4b presents the transport property of C60@SWNTs. The typical p-type characteristic is observed, but the threshold voltage (V_{th}) for hole conductance is found to shift from -28 V to +20 V compared with that of pristine SWNTs, indicating that the p-type behavior of the SWNTs is maintained by the C₆₀ encapsulation. In contrast, the transport property of C₅₉N@SWNTs drastically changes to an n-type semiconductor [Figure 4c]. This n-type characteristic is attributed to the charge transfer between C₅₉N and local parts of SWNTs, suggesting that C₅₉N exert a strong electron donor effect on SWNTs. In addition, it should be noted that the above fullerene-induced characteristics have been observed in many independent SWNTs devices, and they have good reproducibility under measurements performed with different source-drain voltages.

In summary, we have investigated the electrical transport properties of C₆₀ and C₅₉N encapsulated SWNTs by fabricating them as the channels of FET devices at room temperature. Their measurements indicate that C60@SWNTs exhibit the enhanced p-type characteristics compared with the case of pristine SWNTs, whereas C₅₉N@SWNTs show the n-type behavior. The novel transport properties of these encapsulated SWNTs can be explained by the electron donor behavior of C₅₉N, which can modify the electronic structure of SWNTs.

Acknowledgment. The authors thank Prof. K. Tohji and Mr. K. Motomiya for their assistant in TEM observation. This work was supported by a Grant-in-Aid for Scientific Research from the Ministry of Education, Culture, Sports, Science and Technology, Japan.

Supporting Information Available: Raman spectra of fullerenes (C₅₉N and C₆₀) and fullerene encapsulated SWNTs and TEM images for C59N encapsulated SWNTs with lower magnification in PDF format. This material is available free of charge via the Internet at http:// pubs.acs.org.

References

- (1) Hummelen, J. C.; Knight, B.; Pavlovich, J.; Gonzalez, R.; Wudl, F. Science **1995**, 269, 1554-1556.
- Fulop, F.; Rockenbauer, A.; Simon, F.; Pekker, S.; Korecz, L.; Garaj, S.; Janossy, A. Chem. Phys. Lett. 2001, 334, 233-23
- Reuther, U.; Hirsch, A. Carbon 2000, 38, 1539-1549. (4)
- Kumashiro, R.; Tanigaki, K.; Ohashi, H.; Tagmatarchis, N.; Kato, H.; Shinohara, H.; Akasaka, T.; Kato, K.; Aoyagi, S.; Kimura, S.; Takata, M. *Appl. Phys. Lett.* **2004**, *84*, 2154–2156.
- (5) Andreoni, W.; Gygi, F.; Parrinello, M. Chem. Phys. Lett. 1992, 190, 159-
- (6) Shimada, T.; Ohno, Y.; Okazaki, T.; Sugai, T.; Suenaga, K.; Kishimoto, S.; Mizutani, T.; Inoue, T.; Taniguchi, R.; Fukui, N.; Okubo, H.; Shinihara, H. *Physica E* **2004**, *21*, 1089–1092.
- (7) Abe, S.; Sato, G.; Kaneko, T.; Hirata, T.; Hatakeyama, R.; Yokoo, K.; Ono, S.; Omote, K.; Kasama, Y. Jpn. J. Appl. Phys. 2006, 45, 8340-
- Simon, F.; Kuzmany, H.; Nafradi, B.; Feher, T.; Forro, L.; Fulop, F.; Janossy, A.; Korecz, L.; Rockenbauer, A.; Hauke, F.; Hirsch, A. Phys. Rev. Lett. 2006, 97, 136801-1-4.
- Izumida, T.; Hatakeyama, R.; Neo, Y.; Mimura, H.; Omote, K.; Kasama, Y. Appl. Phys. Lett. **2006**, *89*, 093121-1-3.
- (10) Li, Y. F.; Hatakeyama, R.; Kaneko, T.; Izumida, T.; Okada, T.; Kato, T.
- Appl. Phys. Lett. 2006, 89, 093110-1-3.
 Rao, A. M.; Eklund, P. C.; Hodeau, J.-L.; Marques, L.; Nunez-Regueiro, M. Phys. Rev. B 1997, 55, 4766-4773.

JA0773960